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## The Global Weak Solutions of the Compressible Euler Equation with Spherical Symmetry

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### 1 Introduction

The compressible Euler equation for an isentropic gas in  $\mathbf{R}^n$  is given by

$$(1.1) \quad \begin{aligned} \rho_t + \nabla \cdot (\rho \vec{u}) &= 0, \\ (\rho \vec{u})_t + \nabla \cdot (\rho \vec{u} \otimes \vec{u} + p) &= 0, \end{aligned}$$

with the equation of state

$$(1.2) \quad p = a^2 \rho^\gamma,$$

where density  $\rho$ , velocity  $\vec{u}$  and pressure  $p$  are functions of  $x \in \mathbf{R}^n$  and  $t \geq 0$ , while  $a > 0$  and  $\gamma \geq 1$  are given constants.

For one dimensional case ( $n=1$ ), the Cauchy problem for (1.1) with (1.2) has been studied by many authors. Nishida [10] established the existence of global weak solutions, for the first time, for the case  $\gamma = 1$  with arbitrary initial data, and Nishida and Smoller [11] for  $\gamma \geq 1$  but with small initial data, both using Glimm's method. DiPerna [3] extended the latter result to the case of large initial data, using the theory of compensated compactness under the restriction  $\gamma = 1 + 2/(2m+1)$ ,  $m \geq 2$  integers. Ding et al [1], [2] removed this restriction and established the existence of global weak solutions for  $1 < \gamma \leq 5/3$ .

On the other hand, little is known for the case  $n \geq 2$ . No global solutions have been known to exist, but only local classical solutions ([5], [6], [8] and [9]).

In this paper, we will present global weak solutions first for the case  $n \geq 2$ . We will do this, however, only for the case of spherical symmetry with  $\gamma = 1$ . As will be seen below, our proof does not work without these restrictions.

Thus, we look for solutions of the form

$$(1.3) \quad \rho = \rho(t, |x|), \quad \vec{u} = \frac{x}{|x|} \cdot u(t, |x|).$$

Then, denoting  $r = |x|$ , (1.1) becomes

$$(1.4) \quad \begin{aligned} \rho_t + \frac{1}{r^{n-1}} (r^{n-1} \rho u)_r &= 0, \\ \rho (u_t + u u_r) + p_r &= 0, \end{aligned}$$

This equation has a singularity at  $r=0$ . To avoid the difficulty caused by this singularity, we simply deal with the boundary value problem for (1.4) in the domain  $1 \leq r < \infty$  (the exterior of a sphere) with the boundary condition  $u(t, 1) = 0$ , which is identical, under the assumption (1.3), to the usual boundary condition  $\vec{n} \cdot \vec{u} = 0$  for (1.1) where  $\vec{n}$  is the unit normal to the boundary.

Put  $\tilde{\rho} = r^{n-1} \rho$ . Then we get from (1.4)

$$(1.5) \quad \begin{aligned} \tilde{\rho}_t + (\tilde{\rho} u)_r &= 0, \\ u_t + u u_r + \frac{a^2 \gamma \tilde{\rho}_r}{\tilde{\rho}^{2-\gamma} r^{(n-1)(\gamma-1)}} &= \frac{a^2 \gamma (n-1) \tilde{\rho}^{\gamma-1}}{r^n \cdot r^{(n-1)(\gamma-2)}}. \end{aligned}$$

Introduce the Lagrangean mass coordinates

$$(1.6) \quad \tau = t, \quad \xi = \int_1^r \tilde{\rho}(t, r) dr.$$

Then  $\xi > 0$  as long as  $\tilde{\rho} > 0$  for  $r > 1$ , and (1.5) is reformulated as

$$(1.7) \quad \begin{aligned} \tilde{\rho}_\tau + \tilde{\rho}^2 u_\xi &= 0, \\ u_\tau + \frac{a^2 \gamma \tilde{\rho}_\xi}{\tilde{\rho}^{1-\gamma} r^{2\gamma-2}} &= \frac{a^2 \gamma (n-1) \tilde{\rho}^{\gamma-1}}{r^n \cdot r^{(n-1)(\gamma-2)}}. \end{aligned}$$

Put  $v = 1/\tilde{\rho}$  and note that the inverse transformation to (1.6) is given by

$$(1.8) \quad t = \tau, \quad r = 1 + \int_0^\xi v(\zeta, t) d\zeta.$$

Then after changing  $\tau$  to  $t$  and  $\xi$  to  $x$ , (1.7) is written as

$$(1.9) \quad \begin{aligned} v_t - u_x &= 0, \\ u_t + \left( \frac{a^2}{v^\gamma} \right)_x \cdot \frac{1}{r^{(n-1)(\gamma-1)}} &= \frac{a^2 \gamma (n-1) v^{1-\gamma}}{r^n \cdot r^{(n-1)(\gamma-2)}}, \end{aligned}$$

where  $r$  is now defined by  $r = 1 + \int_0^x v(t, \zeta) d\zeta$ .

Now we restrict ourselves to the case  $\gamma = 1$ . Then (1.7) becomes

$$(1.10) \quad \begin{aligned} v_t - u_x &= 0, \\ u_t + \left( \frac{a^2}{v} \right)_x &= \frac{K}{1 + \int_0^x v(t, \zeta) d\zeta}. \end{aligned}$$

where  $K = a^2(n-1)$ .

Let us consider the initial boundary value problem for (1.10) in  $t \geq 0, x \geq 0$  with the following boundary and initial conditions.

$$(1.11) \quad u(0, x) = u_0(x), \quad v(0, x) = v_0(x), \quad \text{for } x > 0,$$

$$(1.12) \quad u(t, 0) = 0, \quad \text{for } t > 0.$$

Let  $BV(\mathbf{R}_+)$  denote the space of functions of bounded variation on  $\mathbf{R}_+ = (0, \infty)$ . Our main result is as follows.

**Theorem ( Main Result )** *Suppose that  $u_0(x), v_0(x) \in BV(\mathbf{R}_+)$ , and that  $v_0(x) \geq \delta_0 > 0$  for all  $x > 0$  with some positive constant  $\delta_0$ . Then (1.10), (1.11) and (1.12) have a global weak solution which belongs to the class*

$$u, v \in L^\infty(0, T; BV(\mathbf{R}_+)) \cap Lip([0, T]; L^1_{loc}(\mathbf{R}_+))$$

for any  $T > 0$ .

The definition of the weak solution will be given in section 4. This theorem can be proved by following Nishida's argument [10] based on Glimm's

method. Indeed this can be seen from the following two simple observations. First, the homogeneous equation corresponding to (1.10),

$$(1.13) \quad \begin{aligned} v_t - u_x &= 0, \\ u_t + \left( \frac{a^2}{v} \right)_x &= 0, \end{aligned}$$

is just the same equation as solved by Nishida [10] using Glimm's method both on the Cauchy problem and the initial boundary value problem. Note that if  $\gamma > 1$ , the homogeneous equation for (1.9) has a variable coefficient and hence does not coincide with the one dimensional Euler equation.

The second observation is that, as long as  $v \geq 0$ , the right hand side of (1.10),

$$(1.14) \quad \frac{K}{1 + \int_0^x v(t, \zeta) d\zeta},$$

is monotone decreasing in  $x$  and has an a priori estimate

$$(1.15) \quad T.V. \left( \frac{K}{1 + \int_0^x v(t, \zeta) d\zeta} \right) \leq K,$$

independent of  $v$ . The one dimensional inhomogeneous Euler equation has been studied in [12]. However, the conditions imposed therein on the inhomogeneous term are not applicable to our (1.14).

These observations allow us to use Nishida's argument [10] to construct global weak solutions to (1.10), (1.11) and (1.12). More precisely, we will first construct, in section 2, approximate solutions of the form

$$\{\text{solution of Riemann problem for (1.13)}\} + \{\text{nonhomogeneous term}\} \times t.$$

This is the main idea of [12]. Then in section 3, we will estimate the total variation of the approximate solutions. Thanks to (1.15), this can be done with a slight modification of Nishida's argument [10]. In section 4, we will show that there exists a subsequence of approximate solutions which converges strongly in  $L^1_{loc}$  for any finite time interval. Finally, for the sake of completeness, we give in Appendix a detailed proof of two lemmas used in section 3. These lemmas are due to Nishida [10], but their proofs are not found in the literature.

## 2 The Difference Scheme

To construct the approximate solutions, we shall use the difference scheme developed in [10]. For  $l, h > 0$ , define

$$(2.1) \quad \begin{aligned} Y &= \{ (n, m); n = 1, 2, 3, \dots, m = 1, 3, 5, \dots \}, \\ A &= \prod_{(m,n) \in Y} [\{nh\} \times ((m-1)l, (m+1)l)] , \end{aligned}$$

where  $l/h$  will be determined later. Choose a point  $\{a_{nm}\} \in A$  randomly, and write  $a_{nm} = (nh, c_{nm})$ . For  $n = 0$ , we put  $c_{0m} = ml$ . We denote approximate solutions by  $u^l$  and  $v^l$ . Mesh lengths  $l$  and  $h$  are chosen so that  $l/h > a/(\inf v^l)$ , for any given  $T > 0$ . We shall show later that there exists a  $\delta > 0$  such that  $\inf v^l \geq \delta > 0$ .

For  $0 \leq t < h$ ,  $ml \leq x < (m+2)l$ ,  $m$  : odd, we define

$$(2.2) \quad \begin{aligned} u^l(t, x) &= u_0^l(t, x) + U^l(t, x)t, \\ v^l(t, x) &= v_0^l(t, x), \end{aligned}$$

where  $u_0^l$  and  $v_0^l$  are the solutions of

$$(2.3) \quad \begin{aligned} v_t - u_x &= 0, \\ u_t + \left( \frac{a^2}{v} \right)_x &= 0, \end{aligned}$$

with initial data

$$(2.4) \quad \begin{aligned} u_0^l(0, x) &= \begin{cases} u_0(ml), & x < (m+1)l, \\ u_0((m+2)l), & x > (m+1)l, \end{cases} \\ v_0^l(0, x) &= \begin{cases} v_0(ml), & x < (m+1)l, \\ v_0((m+2)l), & x > (m+1)l, \end{cases} \end{aligned}$$

and

$$(2.5) \quad U^l(t, x) = \frac{K}{1 + \sum_{j=1}^{\frac{m+1}{2}} v_0((2j-1)l) \cdot 2l} .$$

For  $0 \leq t < h$ ,  $0 \leq x < l$ , we define  $u^l$  and  $v^l$  by (2.2) where  $u_0^l$  and  $v_0^l$  are the solutions of (2.3) with initial boundary data

$$(2.6) \quad u_0^l(0, x) = u_0(l), \quad v_0^l(0, x) = v_0(l), \quad x > 0,$$

$$(2.7) \quad u(t, 0) = 0, \quad t > 0,$$

and

$$(2.8) \quad U^l(t, x) = K.$$

Suppose that  $u^l$  and  $v^l$  are defined for  $0 \leq t < nh$ . For  $nh \leq t < (n+1)h$ ,  $ml \leq x < (m+2)l$ ,  $m$  : odd, we define

$$(2.9) \quad \begin{aligned} u^l(t, x) &= u_0^l(t, x) + U^l(t, x) \cdot (t - nh), \\ v^l(t, x) &= v_0^l(t, x), \end{aligned}$$

where  $u_0^l$  and  $v_0^l$  are the solutions of (2.3) with initial data ( $t=nh$ )

$$(2.10) \quad \begin{aligned} u_0^l(nh, x) &= \begin{cases} u^l(nh - 0, c_{nm}), & x < (m+1)l, \\ u^l(nh - 0, c_{n, m+2}), & x > (m+1)l, \end{cases} \\ v_0^l(nh, x) &= \begin{cases} v^l(nh - 0, c_{nm}), & x < (m+1)l, \\ v^l(nh - 0, c_{n, m+2}), & x > (m+1)l, \end{cases} \end{aligned}$$

and

$$(2.11) \quad U^l(t, x) = \frac{K}{1 + \sum_{j=1}^{\frac{m+1}{2}} v^l(nh - 0, c_{n, 2j-1}) \cdot 2l}.$$

For  $nh \leq t < (n+1)h$ ,  $0 \leq x < l$ , we define  $u^l$  and  $v^l$  as (2.9) where  $u_0^l$  and  $v_0^l$  are the solutions of (2.3) with initial ( $t=nh$ ) boundary data

$$(2.12) \quad u_0^l(nh, x) = u^l(nh - 0, c_{n1}), \quad v_0^l(nh, x) = v^l(nh - 0, c_{n1}), \quad x > 0,$$

$$(2.13) \quad u(t, 0) = 0, \quad t > nh,$$

and  $U^l(t, x)$  is as (2.8).

### 3 Bounds for Approximate Solutions

System (1.6) is hyperbolic provided  $v > 0$ , with the characteristic roots and Riemann invariants given by

$$(3.1) \quad \begin{aligned} \lambda &= -\frac{a}{v}, & r &= u + a \log v, \\ \mu &= \frac{a}{v}, & s &= u - a \log v. \end{aligned}$$

It is well-known, [10], that all shock wave curves in the  $(r,s)$ -plane have the same figure. ( See Figure 1.) The 1-shock wave curve  $S_1$ , starting from  $(r_0, s_0)$  can be expressed in the form

$$(3.2) \quad s - s_0 = f(r - r_0) \quad \text{for } r \leq r_0,$$

and the 2-shock wave curve  $S_2$  can also be expressed in the form

$$(3.3) \quad r - r_0 = f(s - s_0) \quad \text{for } s \leq s_0,$$

where

$$0 \leq f'(x) < 1, \quad f''(x) \leq 0, \quad \lim_{x \rightarrow -\infty} f'(x) = 1.$$

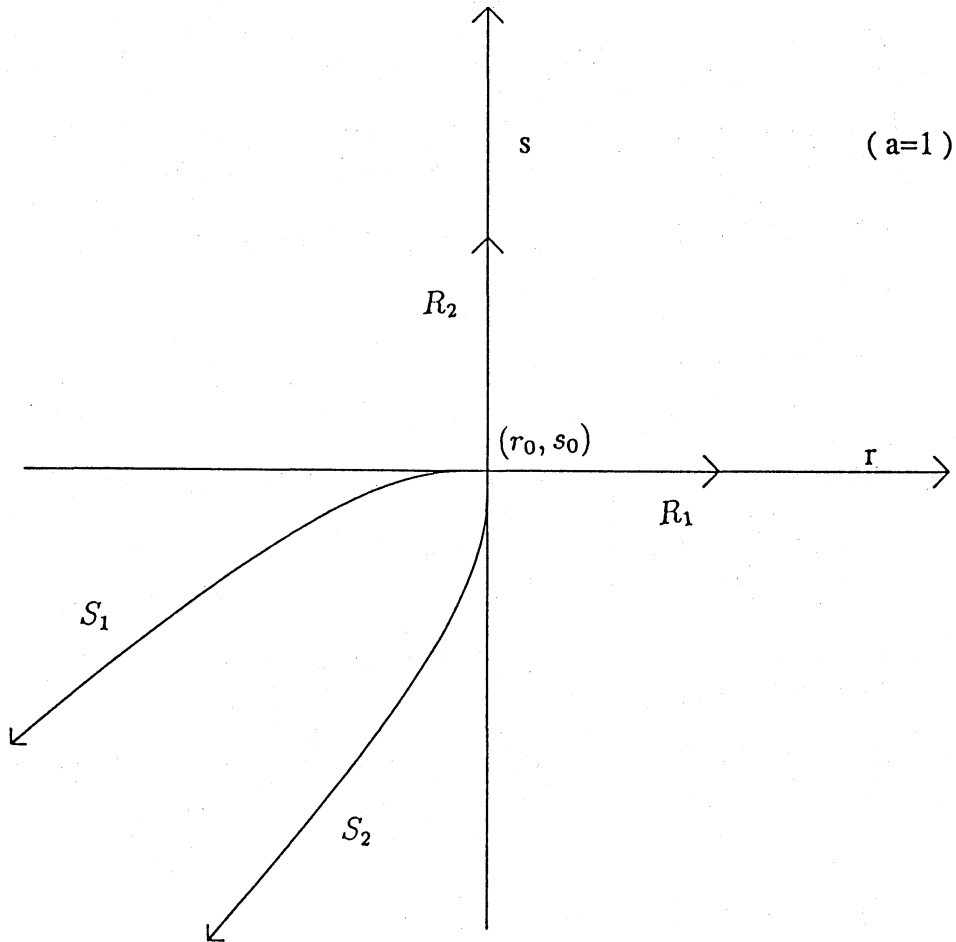


Figure.1



The 1-rarefaction wave curve  $R_1$  can be expressed in the form

$$(3.4) \quad s - s_0 = 0 \quad \text{for } r \geq r_0,$$

and the corresponding expression for the 2-rarefaction wave curve  $R_2$  is

$$(3.5) \quad r - r_0 = 0 \quad \text{for } s \geq s_0.$$

Now we must prepare some lemmas to estimate Riemann invariants. First, let us consider (2.3) with following initial data

$$(3.6) \quad u_0(x) = \begin{cases} u_l, & x < 0, \\ u_r, & x > 0. \end{cases} \quad v_0(x) = \begin{cases} v_l, & x < 0, \\ v_r, & x > 0. \end{cases}$$

**Lemma 3.1** *Let  $u$  and  $v$  are the solutions of (2.3) and (3.6). Then,*

$$(3.7) \quad \begin{cases} r(t, x) \equiv r(u(t, x), v(t, x)) \geq r_0 \equiv \min(r(u_r, v_r), r(u_l, v_l)), \\ s(t, x) \equiv s(u(t, x), v(t, x)) \leq s_0 \equiv \max(s(u_r, v_r), s(u_l, v_l)). \end{cases}$$

Next consider (2.3) in  $t \geq 0, x \geq 0$  with following initial and boundary conditions

$$(3.8) \quad u(0, x) = u_0^+, \quad v(0, x) = v_0^+, \quad \text{for } x > 0,$$

$$(3.9) \quad u(t, 0) = 0, \quad \text{for } t > 0.$$

**Lemma 3.2** *Let  $u$  and  $v$  are the solutions of (2.3), (3.8) and (3.9). Then,*

$$(3.10) \quad \begin{cases} r(t, x) \equiv r(u(t, x), s(t, x)) \geq r(u_0^+, v_0^+), \\ s(t, x) \equiv s(u(t, x), s(t, x)) \leq \max(-r(u_0^+, v_0^+), s(u_0^+, v_0^+)). \end{cases}$$

The above two lemmas were proved in [10]. Using these two lemmas, we can get the following lemma.

**Lemma 3.3** *Let  $u^l$  and  $v^l$  be the approximate solutions defined in section 2 and put  $r_0 = \min r(u_0(x), v_0(x))$  and  $s_0 = \max s(u_0(x), v_0(x))$ . Then, for  $0 < t < T$ ,*

$$(3.11) \quad \begin{cases} r^l(t, x) \equiv r(u^l(t, x), s^l(t, x)) \geq r_0, \\ s^l(t, x) \equiv s(u^l(t, x), s^l(t, x)) \leq \max(-r_0, s_0) + KT \end{cases}$$

Let us consider Riemann problem (2.3) and (3.6). Denote by  $\Delta r$  (resp  $\Delta s$ ) the absolute value of the variation of the Riemann invariant  $r$  (resp  $s$ ) in the first (resp second) shock wave.

**Definition 3.4** *We denote*

$$P(u_l, v_l, u_r, v_r) = \Delta r + \Delta s.$$

Then we have the following lemma.

**Lemma 3.5**

$$(3.12) \quad P(u_1, v_1, u_3, v_3) \leq P(u_1, v_1, u_2, v_2) + P(u_2, v_2, u_3, v_3),$$

where  $u_1, u_2$  and  $u_3$  are arbitrary constants and  $v_1, v_2$  and  $v_3$  are arbitrary positive constants.

We shall prove Lemma 3.5 in the Appendix A.

Denote by  $i_0^{n\pm}$  the straight line segments joining the points  $(0, (n \pm \frac{1}{2})h)$  and  $a_{1n}$ . Let  $F(i_0^{n\pm})$  be the absolute value of the variation of the Riemann invariants for all shocks on  $i_0^{n\pm}$ . Then we also have the following Lemma.

**Lemma 3.6**

$$(3.13) \quad F(i_0^{n+}) \leq F(i_0^{n-}).$$

This lemma 3.6 will be proved in the Appendix B.

We denote

$$\begin{aligned} Z_1 &= \{l - 0, l + 0, 3l - 0, \dots, (2m - 1)l - 0, (2m - 1)l + 0, \dots\}, \\ Z_2 &= \{2l, 4l, 6l \dots 2ml, \dots\}. \end{aligned}$$

Let  $Z_{(n)} = Z_1 \cup Z_2 \cup \{c_{nm}\}$  and line up the elements  $z_{n,i}$  of  $Z_{(n)}$  so that  $z_{n,i} \leq z_{n,i+1}$ . ( We regard  $(2m - 1)l - 0 < (2m - 1)l + 0$  for  $m$  : integer. )  
Let

$$\begin{aligned} F(nh - 0, u^l, v^l) &= \frac{1}{2} F(i_0^{n-}) \\ &+ \sum_{z_{n,i} \in Z_{(n)}} P(u^l(nh - 0, z_{n,i}), v^l(nh - 0, z_{n,i}), u^l(nh - 0, z_{n,i+1}), v^l(nh - 0, z_{n,i+1})), \end{aligned}$$

$$F(nh+0, u^l, v^l) = \frac{1}{2}F(i_0^{n+}) + \sum_{m: \text{odd}} P(u^l(a_{nm}), v^l(a_{nm}), u^l(a_{nm+2}), v^l(a_{nm+2})).$$

Using Lemma 3.5 and Lemma 3.6, we get

$$(3.14) \quad F((n+1)h+0, u^l, v^l) \leq F((n+1)h-0, u^l, v^l).$$

The following equality is obvious from the definition of  $F$ ,  $u^l$  and  $v^l$ .

$$(3.15) \quad F((n+1)h-0, u_0^l, v_0^l) = F(nh+0, u^l, v^l).$$

We also get

$$\begin{aligned} F((n+1)h-0, u^l, v^l) &= F((n+1)h-0, u_0^l, v_0^l) \\ &+ \sum_{m: \text{odd}} P(u^l((n+1)h-0, ml-0), v^l((n+1)h-0, ml-0), \\ &u^l((n+1)h-0, ml+0), v^l((n+1)h-0, ml+0)). \end{aligned}$$

**Lemma 3.7**

$$(3.16) \quad \begin{aligned} &P(u^l((n+1)h-0, ml-0), v^l((n+1)h-0, ml-0), \\ &u^l((n+1)h-0, ml+0), v^l((n+1)h-0, ml+0)) \\ &\leq 2h \left\{ U^l(nh, (m-1)l) - U^l(nh, (m+1)l) \right\}, m : \text{odd}. \end{aligned}$$

*Proof.* From the definition,

$$\begin{aligned} u^l((n+1)h-0, ml-0) &= u_0^l(nh, ml) + U^l(nh, (m-1)l) \cdot h, \\ u^l((n+1)h-0, ml+0) &= u_0^l(nh, ml) + U^l(nh, (m+1)l) \cdot h, \\ v^l((n+1)h-0, ml-0) &= v^l((n+1)h-0, ml+0) = v_0^l(nh, ml). \end{aligned}$$

Therefore we get

$$(3.17) \quad \begin{aligned} &r^l((n+1)h-0, ml-0) - r^l((n+1)h-0, ml+0) \\ &= s^l((n+1)h-0, ml-0) - s^l((n+1)h-0, ml+0) \cdot l \\ &= h \times \left\{ U^l(nh, (m-1)l) - U^l(nh, (m+1)l) \right\} \geq 0 \end{aligned}$$

Thus the following inequality holds.

$$(3.18) \quad \Delta r, \Delta s \leq h \left\{ U^l(nh, (m-1)l) - U^l(nh, (m+1)l) \right\} \leq \Delta r + \Delta s.$$

From (3.18), we get (3.16).  $\square$

Using Lemma 3.7, we get

$$(3.19) \quad F((n+1)h-0, u^l, v^l) - F((n+1)h-0, u_0^l, v_0^l) \leq 2h \sum_{m: \text{odd}} \{ U^l(nh, (m-1)l) - U^l(nh, (m+1)l) \} \leq 2Kh$$

From (3.14), (3.15) and (3.19), we get

$$(3.20) \quad F((n+1)h+0, u^l, v^l) \leq F(nh+0, u^l, v^l) + 2Kh$$

Thus we obtain the following lemma.

**Lemma 3.8**

$$(3.21) \quad F(nh+0, u^l, v^l) \leq F(+0, u^l, v^l) + 2KT \equiv F_0 + 2KT$$

Denote by  $G(\tau)$  the absolute value of the sum of negative variation of  $r^l$  and  $s^l$  for  $t = \tau$ . Then for  $nh \leq \tau < (n+1)h$ , we get

$$(3.22) \quad \begin{aligned} G(\tau) &\leq G(nh) + 2h \sum_{m: \text{odd}} \{ U^l(nh, (m-1)l) - U^l(nh, (m+1)l) \} \\ &\leq G(nh) + 2Kh. \end{aligned}$$

**Lemma 3.9**

$$(3.23) \quad G(nh) \leq 2F(nh+0, u^l, v^l).$$

*Proof.* Denote by  $\delta s$  ( resp  $\delta r$  ) the absolute value of the Riemann invariant  $s$  ( resp  $r$  ) in the first ( resp second ) shock wave. By (3.2) and (3.3),  $\Delta r + \delta s < 2\Delta r$  on the first shock and  $\delta r + \Delta s < 2\Delta s$  on the second shock. So from (3.17), (3.18) and above arguments, we get (3.23).  $\square$

From (3.23), (3.24) and (3.25), for any  $\tau$  (  $nh \leq \tau < (n+1)h$  ),

$$(3.24) \quad \begin{aligned} G(\tau) &\leq G(nh) + 2Kh \leq 2F(nh+0, u^l, v^l) + 2Kh \\ &\leq 2F_0 + 6KT \equiv M_1. \end{aligned}$$

Now we can establish a priori estimates of  $u^l$  and  $v^l$ . Denote by T.V.u the total variation of  $u$ .

**Theorem 3.10** *For any  $T > 0$ , the variation of  $u^l$  and  $v^l$  is bounded uniformly for  $h$  and  $\{a_{mn}\}$ . Their upper bound and lower bound, especially the positive lower bound of  $v^l$ , are also uniformly bounded.*

*Proof.* Denote by  $T.V^+.u$  ( resp  $T.V^-.u$  ) the absolute value of the positive ( resp negative ) variation of  $u$ . Put  $f^l \equiv 2u^l = r^l + s^l$ . Then  $0 \leq f^l(t, 0) \leq Kh$ . Without loss of generality, we assume that  $u_0(x)$  and  $v_0(x)$  are constant outside a bounded interval. Let

$$(3.25) \quad f^l(t, \infty) = r^l(t, \infty) + s^l(t, \infty) \equiv M_2.$$

Then from the definition,

$$f^l(t, 0) + T.V^+.f^l - T.V^-.f^l = f^l(t, \infty).$$

Since  $T.V^-.f^l(t, \cdot) \leq G(t)$  for any  $t$ , (3.24) yields

$$T.V^+.f^l = f^l(t, \infty) + T.V^-.f^l - f^l(t, 0) \leq M_1 + M_2.$$

Thus we get

$$(3.26) \quad T.V.f^l = T.V.2u^l \leq 2M_1 + M_2.$$

From (3.26), we get

$$|f^l| \leq Kh + 2M_1 + M_2 \leq KT + 2M_1 + M_2 \equiv 2M_3.$$

Therefore we get

$$(3.27) \quad |u_l| \leq M_3.$$

Using Lemma 3.2, we get

$$2a \log v^l = r^l - s^l \geq r_0 - (\max(-r_0, s_0) + KT).$$

Thus we get

$$(3.28) \quad v^l \geq \exp \frac{r_0 - (\max(-r_0, s_0) + KT)}{2a} \equiv \frac{1}{M_5}.$$

From the definition,

$$r^l(t, 0) + T.V^+.r^l - T.V^-.r^l = r^l(t, \infty).$$

Using Lemma 3.3 and (3.24),

$$(3.29) \quad T.V^+.r^l = -r^l(0) + T.V^-.r^l + r(t, \infty) \leq -r_0 + M_1 + r(t, \infty).$$

In view of (3.27) and (3.29), there exists a positive constant  $M_6$  such that

$$(3.30) \quad v^l \leq M_6$$

□

**Theorem 3.11** *For any interval  $[x_1, x_2] \subset [0, \infty)$ , we get*

$$(3.31) \quad \int_{x_1}^{x_2} |u^l(t_2, x) - u^l(t_1, x)| + |v^l(t_2, x) - v^l(t_1, x)| dx \\ \leq M \cdot (|t_2 - t_1| + h), \quad 0 \leq t_1, t_2 < T,$$

where  $M$  depends on  $T$ ,  $x_1$  and  $x_2$ , but not on  $l$  and  $h$ .

*Proof.* Without loss of generality, we assume that

$$nh \leq t_1 < (n+1)h < \dots < (n+k)h \leq t_2 < (n+k+1)h.$$

Let

$$\int_{x_1}^{x_2} |u^l(t_2, x) - u^l(t_1, x)| dx \\ \leq I_1 + I_2 + \int_{x_1}^{x_2} |u^l(t_2, x) - u^l((n+k)h + 0, x)| + |u^l(t_1, x) - u^l((n+1)h - 0, x)| dx$$

where

$$I_1 = \int_{x_1}^{x_2} \sum_{i=1}^k |u^l((n+i)h + 0, x) - u^l((n+i)h - 0, x)| dx$$

$$I_2 = \int_{x_1}^{x_2} \sum_{i=1}^{k-1} |u^l((n+i+1)h - 0, x) - u^l((n+i)h + 0, x)| dx$$

and

$$k = \left\lceil \frac{t_2 - t_1}{h} \right\rceil$$

Denote by  $1_{[\alpha, \beta]}$  the characteristic function of the interval  $[\alpha, \beta]$ . We regard  $T.V._{-l < x < l} = T.V._{0 < x < l}$ . Then,

$$\begin{aligned} I_1 &\leq \sum_{i=0}^{k+1} \sum_{m: \text{integer}} \int_{x_1}^{x_2} T.V._{2ml < x < (2m+2)l} u^l((n+i)h - 0, x) \cdot 1_{[2ml, (2m+2)l]} dx, \\ &\leq \left( \left\lceil \frac{t_2 - t_1}{h} \right\rceil + 2 \right) \cdot \left( \sup_{0 \leq t \leq T} T.V. u^l(t, \cdot) \right) \cdot 2l. \end{aligned}$$

$$\begin{aligned} I_2 &\leq \sum_{i=0}^k \sum_m \int_{x_1}^{x_2} \left( T.V._{(2m-1)l < x < (2m+1)l} u_0^l((n+i+1)h - 0, x) \cdot 1_{[(2m-1)l, (2m+1)l]} + Kh \right) dx, \\ &\leq \sum_{i=0}^k 2l \cdot T.V. u_0^l((n+i+1)h - 0, \cdot) + K(x_2 - x_1)h, \\ &\leq \left( \left\lceil \frac{t_2 - t_1}{h} \right\rceil + 1 \right) \cdot \left( 2l \sup_{0 \leq t \leq T} T.V. u_0^l(t, \cdot) + K(x_2 - x_1)h \right). \end{aligned}$$

The remaining terms can be evaluated similarly. For

$$\int_{x_1}^{x_2} |v^l(t_2, x) - v^l(t_1, x)| dx,$$

we also have a similar estimate. Combining these results gives (3.31).  $\square$

#### 4 Convergence of The Approximate Solution

Let  $h_n = T/n$  and  $h_n/l_n = \tilde{\delta} < \delta \equiv 1/M_5$ . Consider the sequence  $(u^{l_n}, v^{l_n})$  ( $n = 1, 2, \dots$ ). Then from Theorem 3.9 and Theorem 3.10, there exists a subsequence which converges in  $L^1_{loc}$  to functions  $(u, v)$  uniformly for  $t \in [0, T]$ . Now we shall prove that  $u(x, t)$  and  $v(x, t)$  are the weak solutions of initial boundary value problem (1.6), (1.7) and (1.8) provided  $\{a_{nm}\}$  is suitably chosen, namely, they satisfy the integral identity

$$\begin{aligned} (4.1) \quad &\int_0^T \int_0^\infty u \phi_t + \left( \frac{a^2}{v} \right) \phi_x + \frac{K}{1 + \int_0^x v(t, \zeta) d\zeta} \cdot \phi dx dt \\ &+ \int_0^\infty u_0(x) \phi(0, x) dx = 0, \end{aligned}$$

$$(4.2) \quad \int_0^T \int_0^\infty v \psi_t - u \psi_x dx dt + \int_0^\infty v_0(x) \psi(0, x) dx = 0,$$

for any smooth functions  $\phi$  and  $\psi$  with compact support in the region  $\{(t, x) : 0 \leq t < T, 0 \leq x < \infty\}$  and  $\phi(t, 0) = 0$ . Now we know that  $u_0^l$  and  $v_0^l$  are weak solutions in each time strip  $nh \leq t < (n+1)h$  so that for each test function  $\phi$  satisfying  $\phi(t, 0) = 0$ ,

$$(4.3) \quad \begin{aligned} & \int_{nh}^{(n+1)h} \int_0^\infty u^l \phi_t + \left( \frac{a^2}{v^l} \right) \phi_x + U^l(t, x) \cdot \phi \, dx dt \\ & + \int_0^\infty u^l(nh + 0, x) \phi(nh, x) \\ & - \int_0^\infty u^l((n+1)h - 0, x) \phi((n+1)h, x) dx = 0 \end{aligned}$$

If we sum this over  $n$ , we get

$$(4.4) \quad \begin{aligned} & \int_0^T \int_0^\infty u^l \phi_t + \left( \frac{a^2}{v^l} \right) \phi_x + U^l(t, x) \cdot \phi \, dx dt + \int_0^\infty u^l(0, x) \phi(0, x) \\ & = - \sum_{k=1}^N \int_0^\infty \{ u^l(kh + 0, x) - u^l(kh - 0, x) \} \cdot \phi(kh, x) dx \end{aligned}$$

where  $N=T/h$ . When  $N \rightarrow \infty$ , the right-hand side of the above equality tends to 0 for almost every  $\{a_{nm}\} \in A$  ( see [4] ). It is immediate to see that

$$\int_0^\infty u^l(0, x) \phi(0, x) dx \rightarrow \int_0^\infty u_0(x) \phi(0, x) dx \quad (N \rightarrow \infty).$$

#### Lemma 4.1

$$(4.5) \quad U^l(t, x) \rightarrow \frac{K}{1 + \int_0^x v(t, \zeta) d\zeta} \quad (N \rightarrow \infty).$$

locally uniformly for  $t$  and  $x$ .

*Proof.* Let  $nh \leq t < (n+1)h$ ,  $x \in ((m-1)l, (m+1)l)$ ,  $m : \text{odd}$ . Then

$$(4.6) \quad \left| \int_0^x v^l(nh, \zeta) d\zeta - \sum_{j=1}^{\frac{m+1}{2}} v^l(nh, c_{2j-1n}) \right| \leq \|v^l\|_\infty \cdot l.$$

On the other hand

$$(4.7) \quad \int_0^x v^l(t, \zeta) d\zeta \rightarrow \int_0^x v(t, \zeta) d\zeta \quad (N \rightarrow \infty).$$



locally uniformly for  $t$  and  $x$ .

We get

$$\begin{aligned}
 (4.8) \quad & \left| \int_0^x v^l(t, \zeta) d\zeta - \int_0^x v^l(nh, \zeta) d\zeta \right| \\
 & \leq \int_0^x \sum_{m: \text{odd}} T \cdot V \cdot (m-1)l < \zeta < (m+1)l v^l(nh, \cdot) \cdot 1_{[(m-1)l, (m+1)l]} d\zeta \\
 & \leq \sup_{0 \leq t \leq T} T \cdot V \cdot v^l \cdot 2l.
 \end{aligned}$$

From (4.6), (4.7) and (4.8), we get (4.5).  $\square$

For each test function  $\psi$ ,  $v^l$  also satisfies,

$$\begin{aligned}
 (4.9) \quad & \int_0^T \int_0^\infty (v^l \psi_t - u^l \psi_x) dx dt + \int_0^\infty v^l(0, x) \psi(0, x) dx \\
 & = - \sum_{k=1}^N \int_0^\infty \{v^l(kh + 0, x) - v^l(kh - 0, x)\} \cdot \psi(kl, x) dx \\
 & \quad - I_1 - I_2.
 \end{aligned}$$

where

$$I_1 = \sum_{n=0}^{N-1} \int_{nh}^{(n+1)h} U^l(t, 0)(t - nh) \psi(t, 0) dt$$

and

$$I_2 = \sum_{n=0}^{N-1} \sum_{m: \text{odd}} \int_{nh}^{(n+1)h} \{U^l(t, ml + 0) - U^l(t, ml - 0)\} (t - nh) \psi(t, ml) dt.$$

The first term of the the right-hand side of equality (4.9) tends to 0 for almost every  $\{a_{nm}\} \in A$  ( see [4]). It is also immediate to see that

$$\int_0^\infty v^l(0, x) \psi(0, x) dx \rightarrow \int_0^\infty v_0(x) \psi(0, x) dx \quad (N \rightarrow \infty).$$

We shall show that  $I_1, I_2 \rightarrow 0$  as  $N \rightarrow \infty$ .

$$\begin{aligned}
 (4.10) \quad & I_1 \leq \| \psi \|_\infty \sum_{n=0}^{N-1} \int_{nh}^{(n+1)h} U^l(t, 0)(t - nh) dt \\
 & \leq \| \psi \|_\infty \sum_{n=0}^{N-1} \int_{nh}^{(n+1)h} K(t - nh) dt \\
 & \leq \| \psi \|_\infty h a^2 T.
 \end{aligned}$$

$$\sum_{m: \text{odd}} \int_{nh}^{(n+1)h} \{U^l(t, ml + 0) - U^l(t, ml - 0)\} (t - nh) \psi(t, ml) dt \leq K \|\psi\|_{\infty} h^2.$$

Thus we get

$$(4.11) \quad I_2 \leq \|\psi\|_{\infty} \sum_{n=0}^{N-1} Kh^2 \leq K \|\psi\|_{\infty} hT$$

From above arguments, we can conclude that  $u$  and  $v$  satisfy (4.1) and (4.2). Thus we obtain our main result.

**Theorem 4.2 ( Main Result )** Suppose that  $u_0(x), v_0(x) \in BV(\mathbf{R}_+)$ , and that  $v_0(x) \geq \delta_0 > 0$  for all  $x > 0$  with some positive constant  $\delta_0$ . Then (1.10), (1.11) and (1.12) have a global weak solution which belongs to the class

$$u, v \in L^{\infty}(0, T; BV(\mathbf{R}_+)) \cap Lip([0, T]; L^1_{loc}(\mathbf{R}_+))$$

for any  $T > 0$ .

## Appendix

### A Proof of Lemma 3.5

Let  $g(x) = -f(-x)$ , and put

$$P(u_1, v_1, u_2, v_2) = \Delta r_1 + \Delta s_1$$

$$P(u_2, v_2, u_3, v_3) = \Delta r_2 + \Delta s_2$$

$$P(u_1, v_1, u_3, v_3) = \Delta r_3 + \Delta s_3$$

Then it is obvious that

$$\begin{aligned} & \Delta r_3 + g(\Delta s_3) + \Delta s_3 + g(\Delta r_3) \\ & \leq \Delta r_1 + \Delta r_2 + \Delta s_1 + \Delta s_2 + g(\Delta r_1) + g(\Delta r_2) + g(\Delta s_1) + g(\Delta s_2) \end{aligned}$$

We notice that  $f'' \leq 0$  and hence

$$\leq \Delta r_1 + \Delta r_2 + \Delta s_1 + \Delta s_2 + g(\Delta r_1 + \Delta r_2) + g(\Delta s_1 + \Delta s_2).$$

Let  $x + g(x) = h(x)$ ,  $\Delta r_3 = p'$ ,  $\Delta s_3 = q'$ ,  $\Delta r_1 + \Delta r_2 = p$  and  $\Delta s_1 + \Delta s_2 = q$ .

Then

$$(A.1) \quad h(p') + h(q') \leq h(p) + h(q).$$

Put  $K = h(p') + h(q')$ . We shall estimate  $p + q$  from below under the restriction (A.1). To do this, as  $h$  is monotone increasing function, we must estimate  $p + q$  from below under the restriction

$$(A.2) \quad h(p) + h(q) = K.$$

We do this by using Lagrange's method of indeterminate coefficients.

Put  $G(p, q, \lambda) = p + q + \lambda (h(p) + h(q) - K)$ . Then

$$G_p = 1 + \lambda h'(p) = 0, \quad G_q = 1 + \lambda h'(q) = 0.$$

Because  $h''(x) > 0$ , we get  $p = q$ . So  $p + q$  attains its extremum at  $p = q$ .

We can show that when  $p = q$ ,  $p + q$  is minimum under the restriction (A2).

Therefore

$$h(p) = h(q) = \frac{K}{2} = \frac{h(p') + h(q')}{2} \geq h\left(\frac{p' + q'}{2}\right).$$

Hence it follows that

$$p = q \geq \frac{p' + q'}{2}.$$

Thus we get

$$(A.3) \quad p + q \geq p' + q'.$$

which proves Lemma 3.5.

## B Proof of Lemma 3.6

To prove Lemma 3.6, we must check the following 12 cases:

- 1)  $c_{1n} < l$ ,
  - (1)  $S_2$  crosses  $i_0^{n-}$ ,
  - (2)  $R_2$  crosses  $i_0^{n-}$ ,
  - (3) no wave cross  $i_0^{n-}$ .
- 2)  $c_{1n} \geq l$ ,
  - (1)  $S_2$  and  $S_1$  cross  $i_0^{n-}$ ,
  - (2)  $R_2$  and  $S_1$  cross  $i_0^{n-}$ ,
  - (3)  $S_2$  and  $R_1$  cross  $i_0^{n-}$ ,
  - (4)  $R_2$  and  $R_1$  cross  $i_0^{n-}$ ,
  - (5)  $S_1$  crosses  $i_0^{n-}$ ,
  - (6)  $R_1$  crosses  $i_0^{n-}$ ,
  - (7)  $S_2$  crosses  $i_0^{n-}$ ,
  - (8)  $R_2$  crosses  $i_0^{n-}$ ,
  - (9) no wave cross  $i_0^{n-}$ .

Put  $r_+^{n-1} = r^l(a_{1n-1})$ ,  $s_+^{n-1} = s^l(a_{1n-1})$ ,  $r_-^{n-1} = -s_-^{n-1} = r^l((n-1)h + 0, 0)$ , and  $\delta_{n-1} = U^l(a_{1n-1})$ .  
 Put  $r_+^{n-1'} = r^l((n-1)h + 0, 2l)$  and  $s_+^{n-1'} = s^l((n-1)h + 0, 2l)$ .  
 Put  $A = (r_-^{n-1}, s_-^{n-1})$ ,  $B = (r_+^{n-1}, s_+^{n-1})$  and  $B' = (r_+^{n-1'}, s_+^{n-1'})$ .  
 Put  $C = (r_+^{n-1} + Kh, s_+^{n-1} + Kh)$ ,  
 ( resp  $= (r_+^{n-1'} + \delta_{n-1}h, s_+^{n-1'} + \delta_{n-1}h)$  ) if  $c_{1n} < l$  (resp  $c_{1n} \geq l$ ).  
 If  $R_2$  crosses  $i_0^{n+}$ ,  $F(i_0^{n+}) = 0 \leq F(i_0^{n-})$ , so that it is sufficient to consider the cases when  $S_2$  crosses  $i_0^{n+}$ .

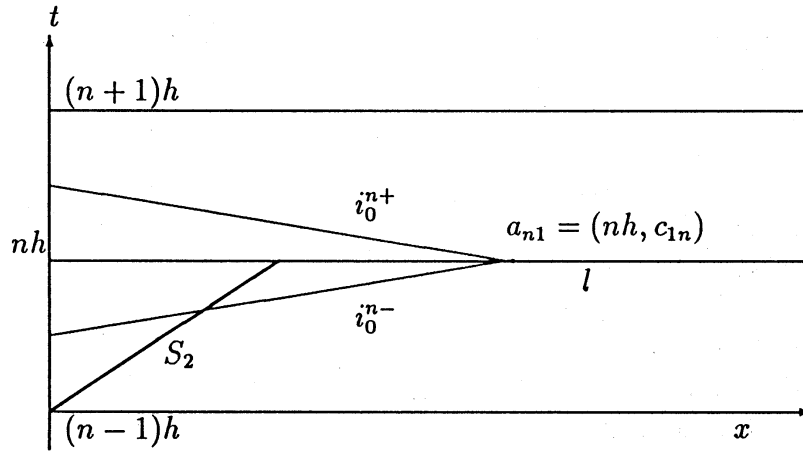


Figure.2

1)  $c_{1n} < l$ .

(1)  $S_2$  crosses  $i_0^{n-}$  ( Figure 2 ). Denote by I ( resp II ) the halfspace  $\{(r, s) | r + s < 0\}$  ( resp  $\{(r, s) | r + s \geq 0\}$  ). )

i)  $C \in I$ .

In this case  $S_2$  crosses  $i_0^{n+}$ . Denote by  $V(PQ)$  the absolute value of the total variation of  $r$  and  $s$  by the line segment  $PQ$ . From Figure.3,

$$F(i_0^{n+}) = V(A'C) \leq V(A'C') = V(AB) = F(i_0^{n-}).$$

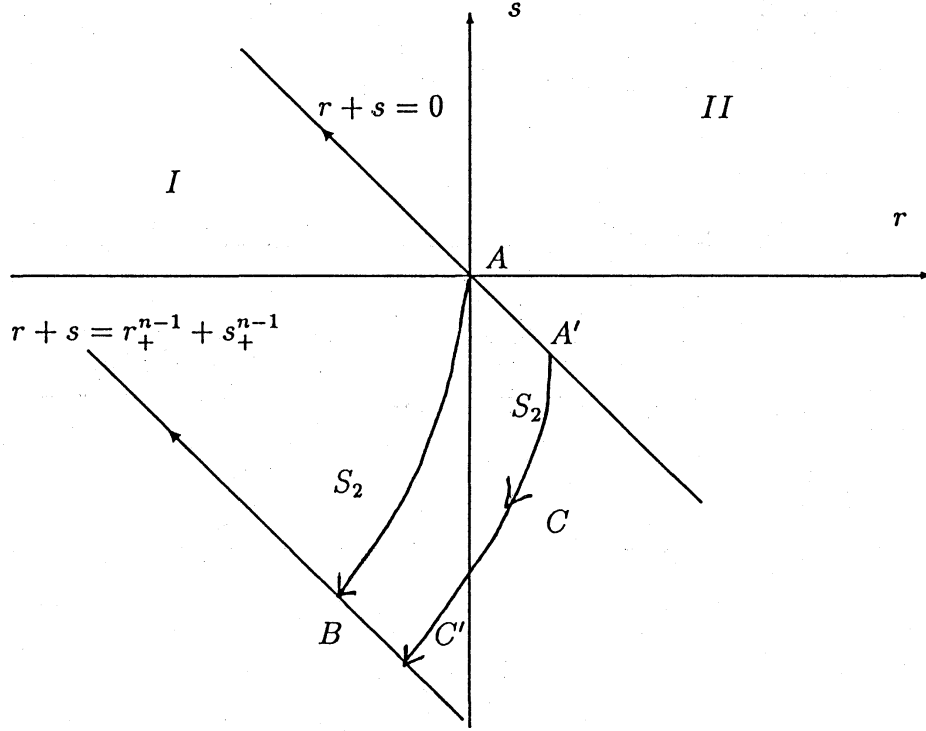


Figure.3

ii)  $C \in II$ .

In this case  $R_2$  crosses  $i_0^{n+}$ . Then

$$(B.1) \quad F(i_0^{n-}) \geq F(i_0^{n+}) = 0.$$

(2)  $R_2$  crosses  $i_0^{n-}$ .

In this case  $B \in II$  so that  $R_2$  crosses  $i_0^{n+}$ . Then

$$(B.2) \quad F(i_0^{n-}) = F(i_0^{n+}) = 0.$$

(3) no wave crosses  $i_0^{n-}$ .

In this case  $(r_+^{n-1}, s_+^{n-1})$  is on the line  $r + s = 0$ . Hence  $C \in II$ . It is obvious that (B.3) also holds.

2)  $c_{1n} \geq l$ .

(1)  $S_2$  and  $S_1$  cross  $i_0^{n-}$ . ( Figure.4 )

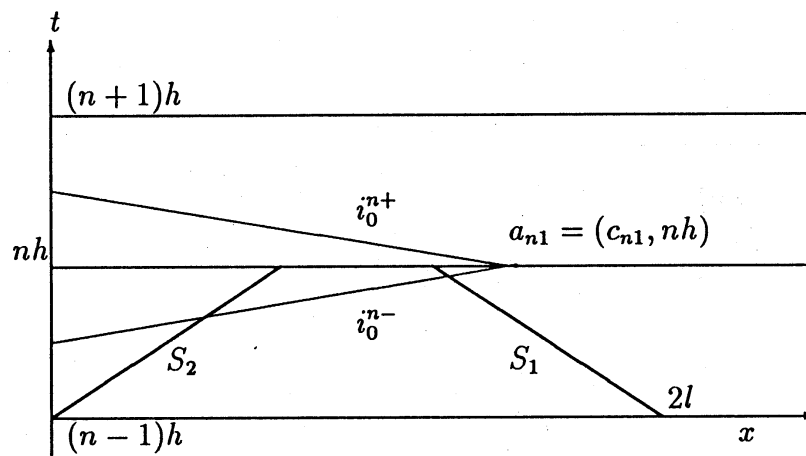


Figure.4

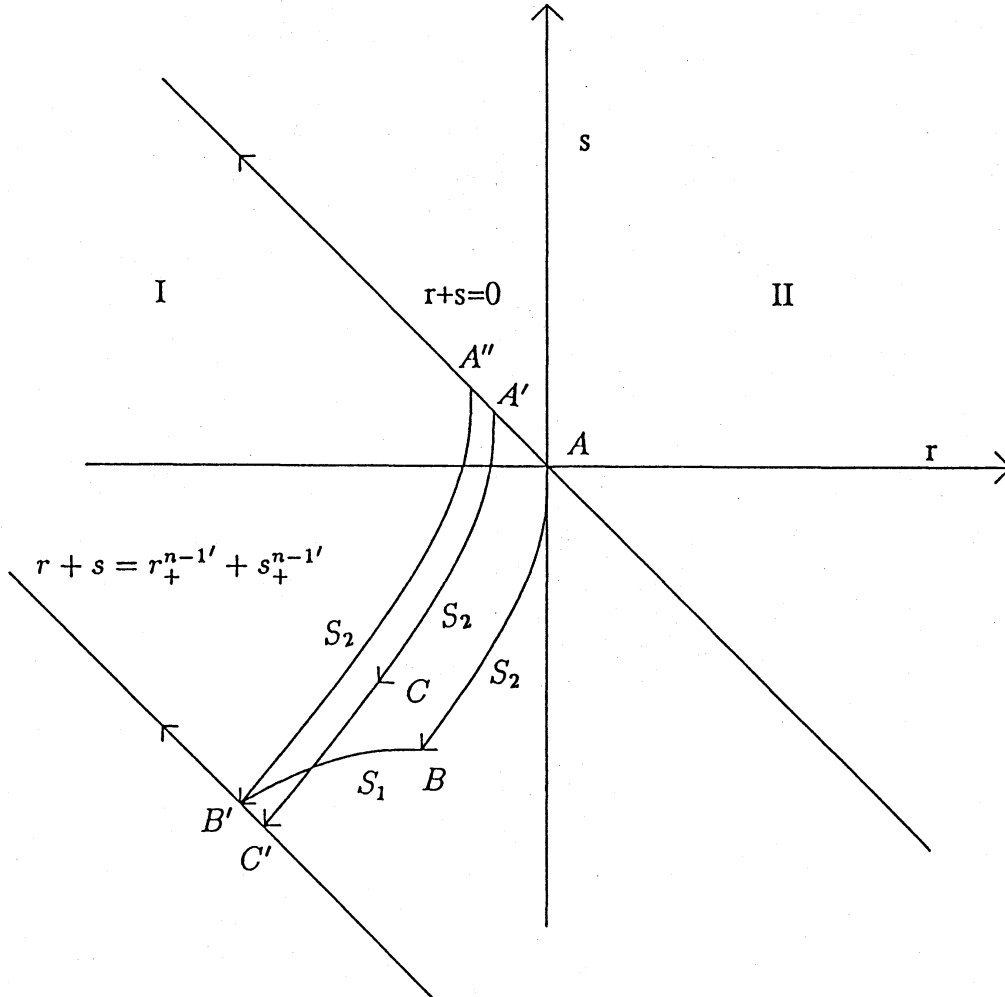


Figure.5

i)  $C \in I$ .

From Figure.5,

$$F(i_0^{n+}) = V(A'C) \leq V(A'C') = V(A''B') = V(AB') = F(i_0^{n-}).$$

ii)  $C \in II$  implies that  $R_2$  crosses  $i_0^{n+}$ . So we get (B2).



(2)  $R_2$  and  $S_1$  cross  $i_0^{n-}$ .

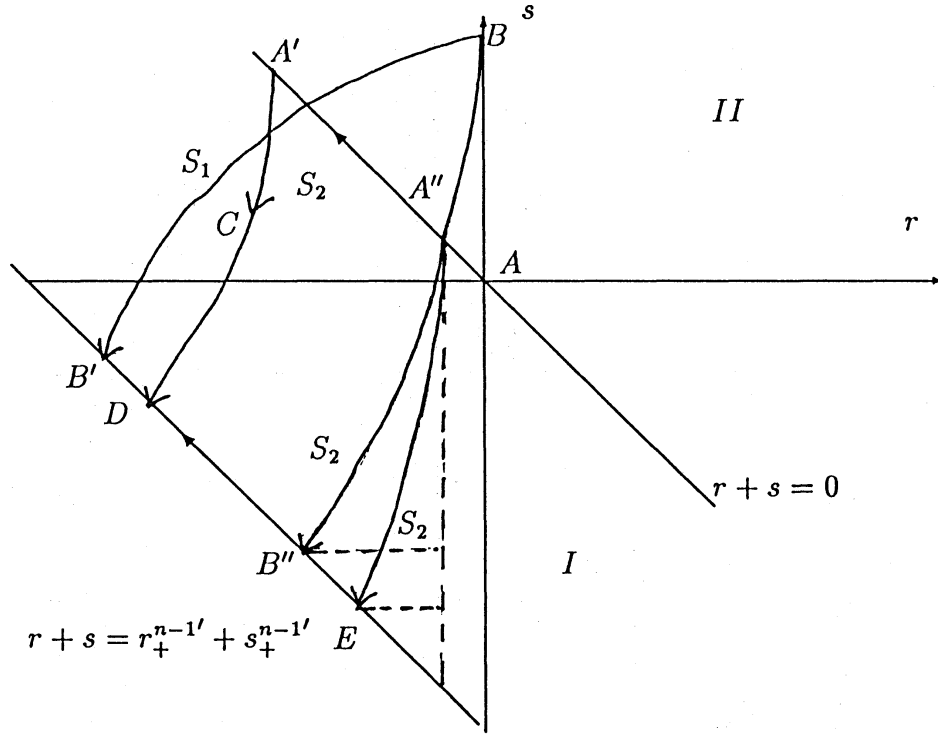


Figure.6

i)  $C \in I$ .

From Figure.6,

$$\begin{aligned} F(i_0^{n+}) &= V(A'C) \leq V(A'D) = V(A''E) = V(A''B'') \\ &\leq V(BB'') = V(BB') = F(i_0^{n+}) \end{aligned}$$

ii)  $C \in II$ .

Thus  $R_2$  crosses  $i_0^{n+}$ , and we get (B2).

(3)  $S_2$  and  $R_1$  cross  $i_0^{n-}$ .

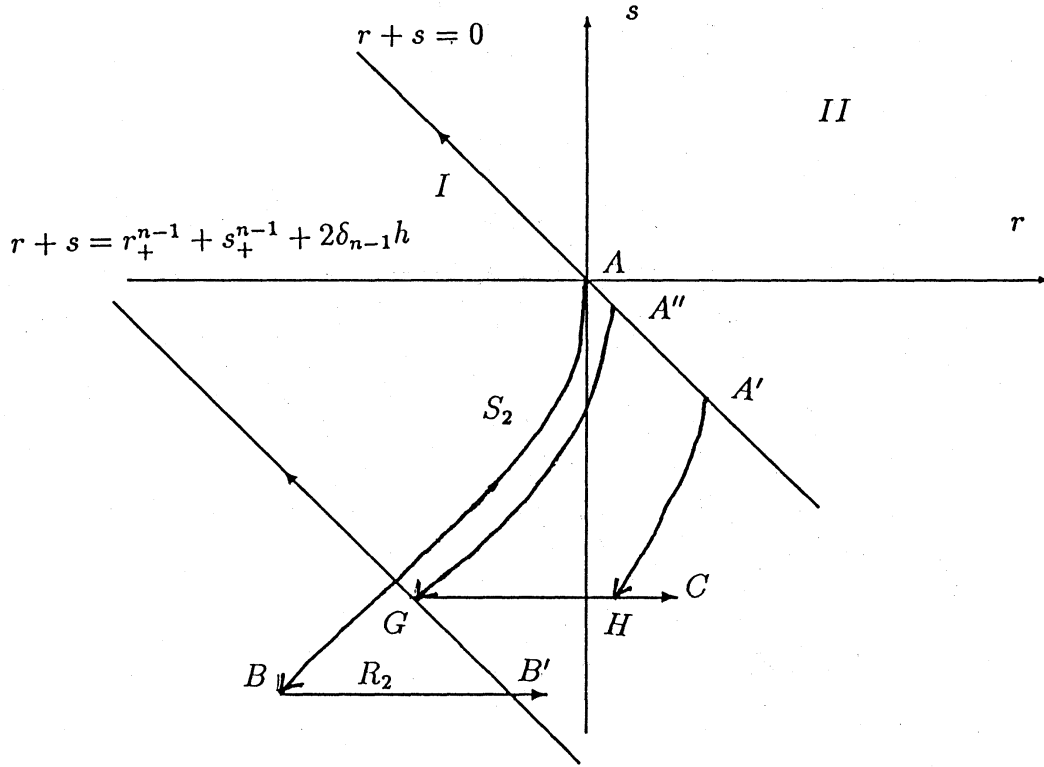


Figure.7

Put  $G = (r_+^{n-1} + \delta_{n-1}h, s_+^{n-1} + \delta_{n-1}h)$  and  $H = (r^l(a_{1n}), s^l(a_{1n}))$ . Then  $H$  is on the line  $CG$ .

i)  $H \in I$ .

From Figure.7,

$$F(i_0^{n+}) = V(A'H) \leq V(A''G) \leq V(AB) = F(i_0^{n-}).$$

ii)  $H \in II$ , so

$R_2$  crosses  $i_0^{n+}$ , and we get (B2).

(4)  $R_2$  and  $R_1$  cross  $i_0^{n-}$ .

In this case,  $R_2$  crosses  $i_0^{n+}$ . So we get (B3).

(5)  $S_1$  crosses  $i_0^{n-}$ .

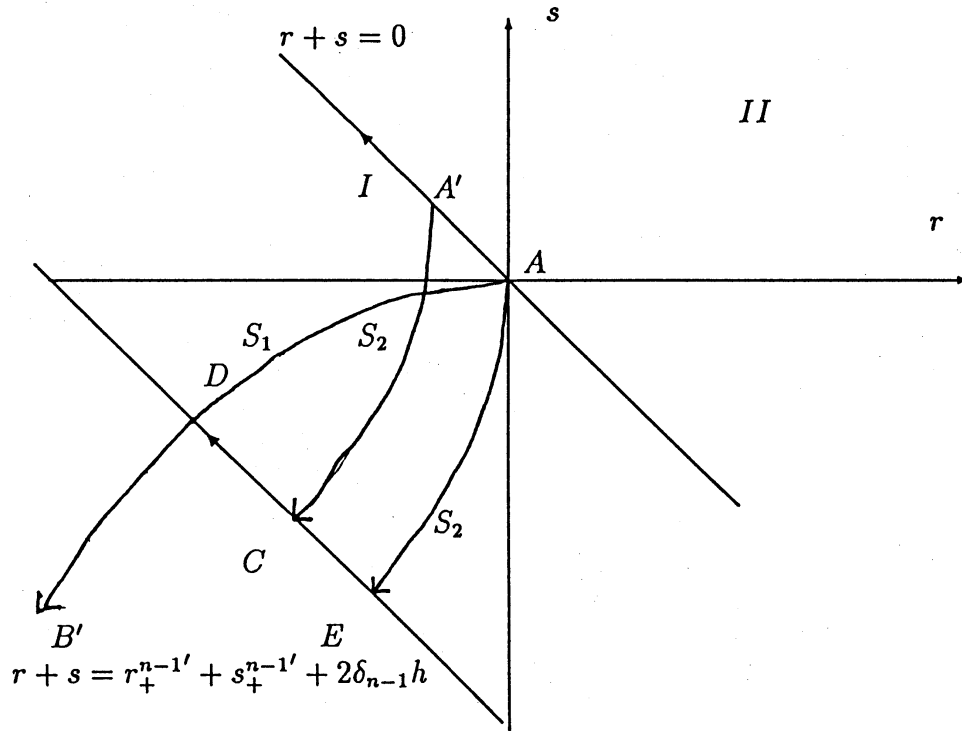


Figure.8

i)  $C \in I$ .

From Figure.8,

$$\begin{aligned} F(i_0^{n+}) &= V(A'C) = V(AE) = V(AD) \\ &\leq V(AB') = F(i_0^{n-}) \end{aligned}$$

Thus we get (B1).

ii)  $C \in II$ .

$R_2$  crosses  $i_0^{n+}$ . So we get (B2).

(6)  $R_1$  crosses  $i_0^{n-}$ .

In this case, it is obvious that  $F(i_0^{n+}) = 0$ . Hence we get (B3).

Cases (7), (8) and (9) are almost the same as cases (1), (2) and (3) in 1). Thus, we obtain Lemma 3.6.

## References

- [1] X. Ding, G. Chen and P. Luo, *Convergence of the Lax-Friedrichs scheme for isentropic gas dynamics*, (I), (II), Acta Math. Sci., 5, (1985), 483-500, 501-540.
- [2] X. Ding, G. Chen and P. Luo, *Convergence of the Lax-Friedrichs scheme for isentropic gas dynamics*, (III), Acta Math. Sci., 6, (1986), 75-120.
- [3] R. DiPerna, *Convergence of the viscosity method for isentropic gas dynamics*, Commun. Math. Phys., 91, (1983), 1-30.
- [4] J. Glimm, *Solutions in the large for nonlinear hyperbolic systems of equations*, Comm. Pure Appl. Math., 18, (1965), 697-715.
- [5] T. Kato, *The Cauchy problem for quasi-linear symmetric hyperbolic systems*, Arch. Rational Mech. Anal., 58, (1975), 181-205.
- [6] P. D. Lax, *Hyperbolic systems of conservation laws and the mathematical theory of shock waves*, SIAM Reg. Conf. Lecture 11, Philadelphia, 1973.
- [7] T. P. Liu and J. Smoller, *On the vacuum state for the isentropic gas dynamics equations*, Advances in Applied Math., 1, (1980), 345-359.
- [8] A. Majda, *Compressible Fluid Flow and Systems of Conservation Laws in Several Space Variables*, Springer-Verlag New York Inc. 1984.
- [9] T. Makino, S. Ukai and S. Kawashima, *Sur la solution à support compact de l'équation d'Euler compressible*, Japan J. Appl. Math., 3, (1986), 249-257.

- [10] T. Nishida, *Global solutions for an initial boundary value problem of a quasilinear hyperbolic system*, Proc. Japan Acad., 44, (1968), 642-646.
- [11] T. Nishida and J. Smoller, *Solutions in the large for some nonlinear hyperbolic conservations*, Comm. Pure Appl. Math., 26, (1973), 183-200.
- [12] L. A. Ying and C. H. Wang, *Global solutions of the Cauchy problem for a nonhomogeneous quasilinear hyperbolic system*, Comm. Pure Appl. Math., 33, (1980), 579-597.